

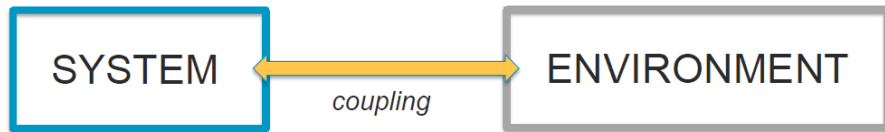
# Arnoldi-Lindblad time evolution: Faster-than-the-clock algorithm for the spectrum of (Floquet) open quantum systems

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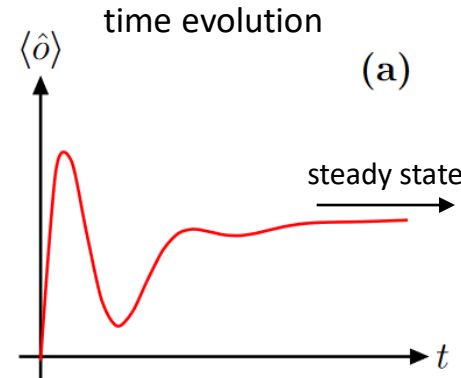
## Lindblad master equation

$$\partial_t \hat{\rho} = \underbrace{-i[\hat{H}, \hat{\rho}]}_{\text{closed system dynamics}} + \underbrace{\sum_j \frac{\gamma_j}{2} (2\hat{L}_j \hat{\rho} \hat{L}_j^\dagger - \{\hat{L}_j^\dagger \hat{L}_j, \hat{\rho}\})}_{\text{coupling with environment}} = \mathcal{L}[\hat{\rho}]$$

dissipation rate
Lindblad operator  
"dissipation operator"

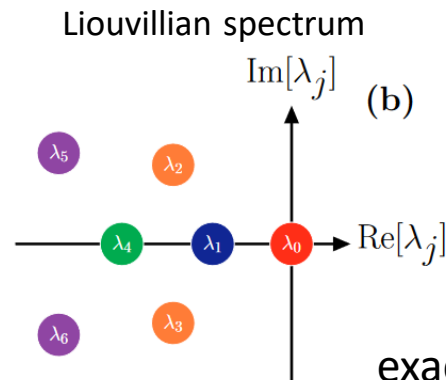
$$\mathcal{L}\hat{\rho}_j = \lambda_j \hat{\rho}_j$$

**Liouvillian superoperator**



main interest in "slow" states  
e.g. the steady state

$$\hat{\rho}(t) = e^{\mathcal{L}t} \hat{\rho}(0) = \sum_j c_j e^{\lambda_j t} \hat{\rho}_j$$



Liouvillian eigendecomposition highly desired

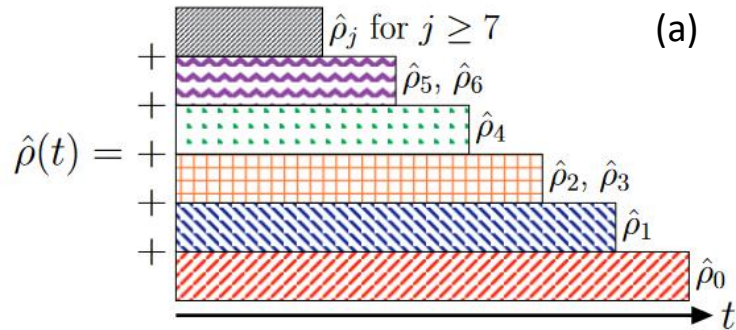


exact diagonalisation often infeasible while time evolution is still accessible

## THE METHOD

### Time evolution

$$\hat{\rho}(t) = \hat{\rho}_{ss} + \sum_{j \geq 1} c_j e^{\lambda_j t} \hat{\rho}_j$$

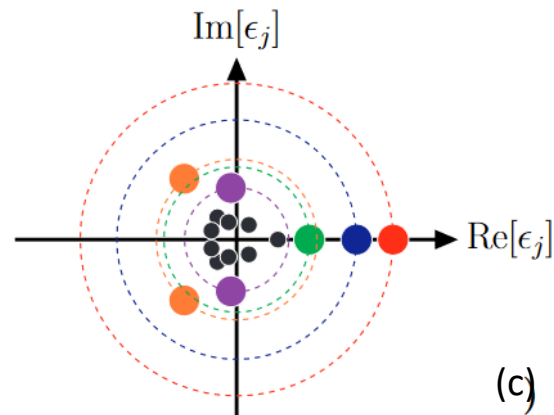
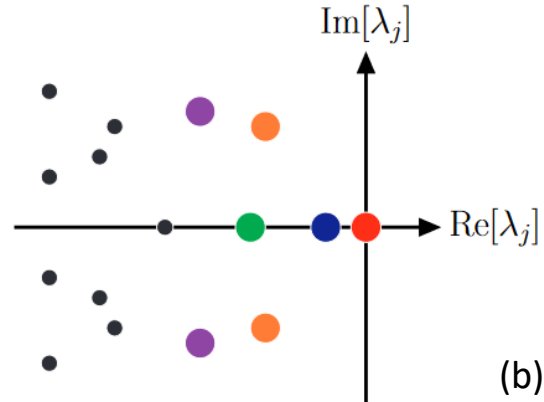


define a time step  $T$

$$\mathcal{E} = e^{\mathcal{L}T}$$

$$\hat{\rho}(t) = [e^{\mathcal{L}T}]^N \hat{\rho}(0) = \mathcal{E}^N \hat{\rho}(0)$$

### Spectrum transformation



$$\mathcal{E} \hat{\rho}_j = \epsilon_j \hat{\rho}_j = e^{\lambda_j T} \hat{\rho}_j \Leftrightarrow \mathcal{L} \hat{\rho}_j = \lambda_j \hat{\rho}_j$$

### Krylov and Arnoldi method

$$K_n = \{\hat{\rho}(0), \hat{\rho}(T), \hat{\rho}(2T), \dots, \hat{\rho}(nT)\}$$

$$= \{\hat{\rho}(0), \mathcal{E} \hat{\rho}(0), \mathcal{E}^2 \hat{\rho}(0), \dots, \mathcal{E}^n \hat{\rho}(0)\}$$

Arnoldi iteration

Gram-Schmidt orthonormalisation



$$\mathcal{E}_n^{eff} = \begin{bmatrix} e_{1,1} & e_{1,2} & e_{1,3} & \dots & e_{1,n} \\ e_{2,1} & e_{2,2} & e_{2,3} & \dots & e_{2,n} \\ 0 & e_{3,2} & e_{3,3} & \dots & e_{3,n} \\ \vdots & \ddots & \ddots & \ddots & \vdots \\ 0 & \dots & 0 & e_{n,n-1} & e_{n,n} \end{bmatrix}$$

check convergence

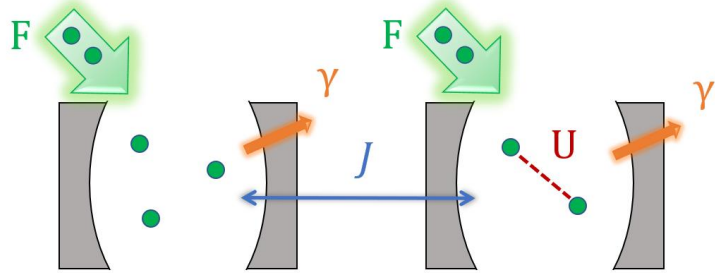
$$\|\mathcal{E} \hat{\rho}_j^{eff} - \epsilon_j^{eff} \hat{\rho}_j^{eff}\| < \tau$$



$$\mathcal{L} \hat{\rho}_j^{eff} = \lambda_j^{eff} \hat{\rho}_j^{eff}$$

# THE DRIVEN DISSIPATIVE BOSE HUBBARD DIMER

$$\hat{H} = \sum_{j=1}^2 \left[ -\Delta \hat{a}_j^\dagger \hat{a}_j + \frac{U}{2} \hat{a}_j^\dagger \hat{a}_j^\dagger \hat{a}_j \hat{a}_j + F(\hat{a}_j^\dagger + \hat{a}_j) \right] - J(\hat{a}_1^\dagger \hat{a}_2 + \hat{a}_2^\dagger \hat{a}_1)$$

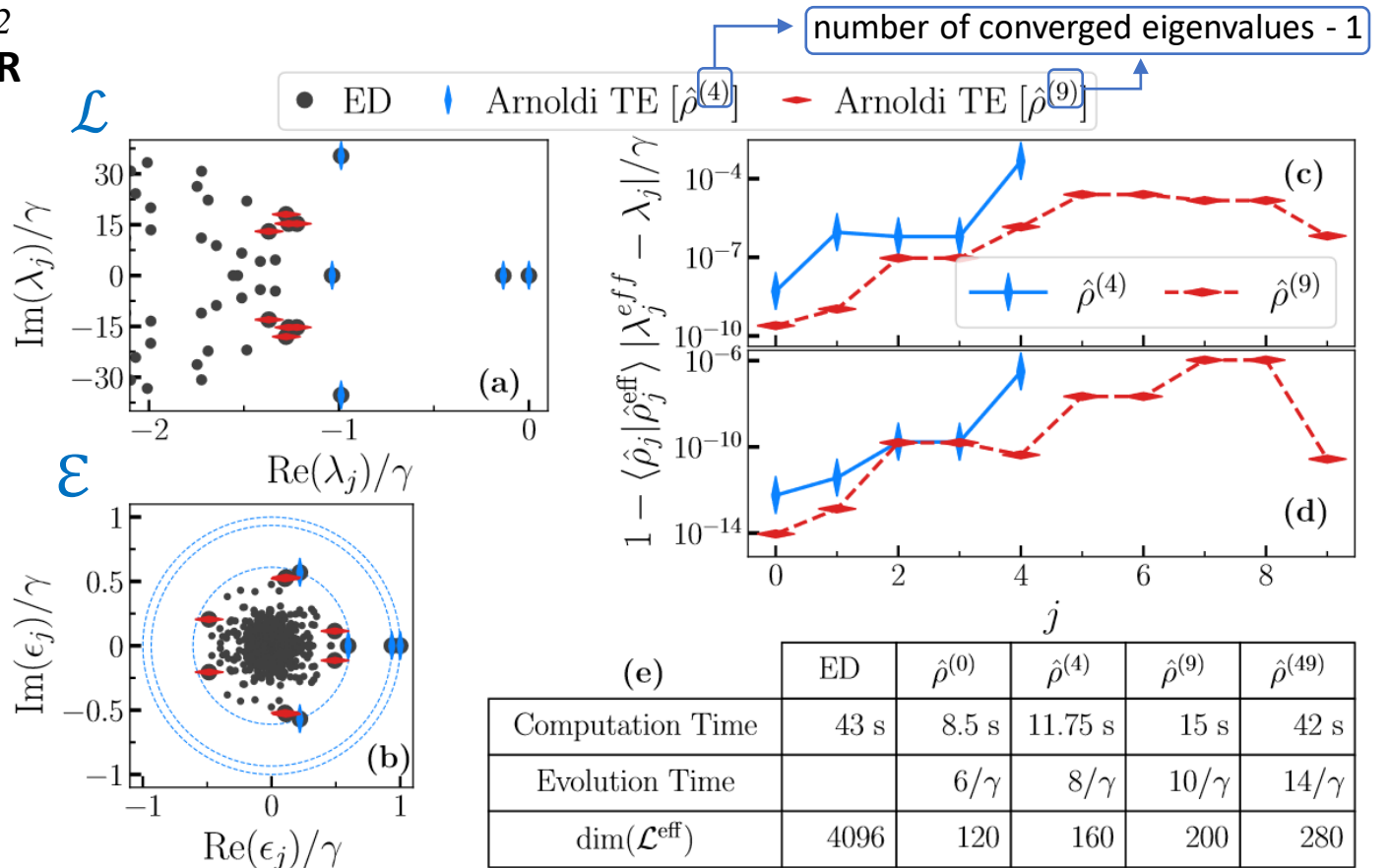


$$\partial_t \hat{\rho} = -i[\hat{H}, \hat{\rho}] + \frac{\gamma}{2} \sum_{j=1}^2 (2\hat{a}_j \hat{\rho} \hat{a}_j^\dagger - \{\hat{a}_j^\dagger \hat{a}_j, \hat{\rho}\})$$

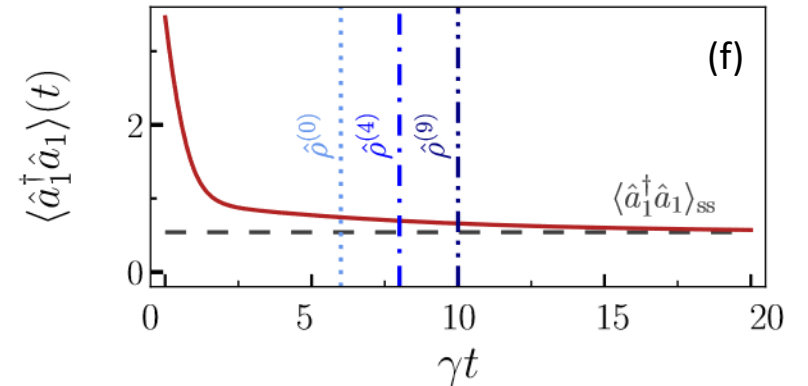
parameters:

$$U = 20\gamma \quad \Delta = 5\gamma \quad n_{max} = 7$$

$$F = 4,5\gamma \quad J = 10\gamma$$



- very small error on eigendecomposition
- faster and more precise



# FLOQUET SYSTEMS

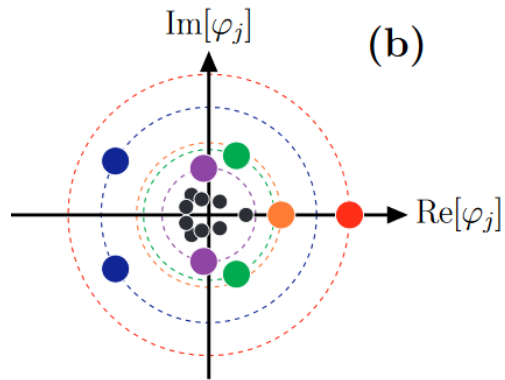
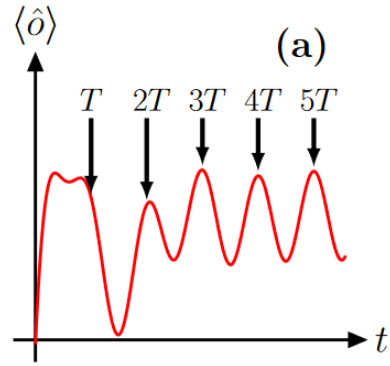
$$\partial_t \hat{\rho}(t) = \mathcal{L}(t) \hat{\rho}(t),$$

$$\mathcal{L}(t+T) = \mathcal{L}(t)$$

$$\hat{\rho}(T) = \mathcal{E} \hat{\rho}(0) \equiv \mathcal{F} \hat{\rho}(0)$$

$$= e^{\mathcal{L}^F T} \hat{\rho}(0)$$

"Floquet map"



stroboscopic eigenspectrum

$$\mathcal{F} \hat{\rho}_j^F = \varphi_j^F \hat{\rho}_j^F$$

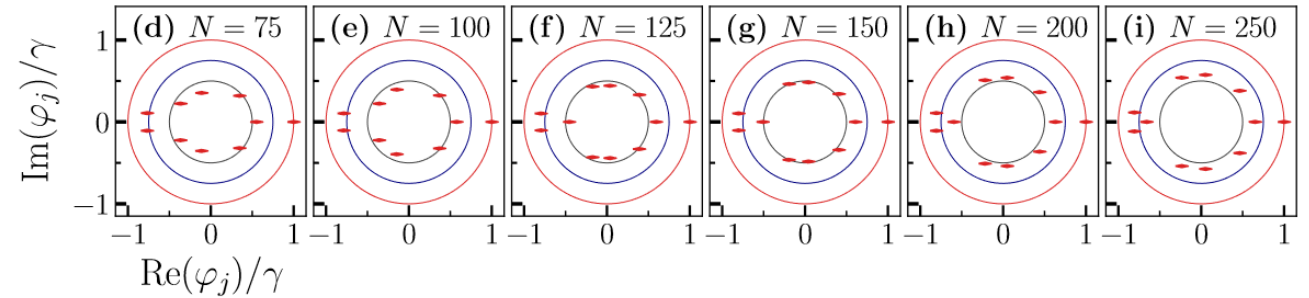
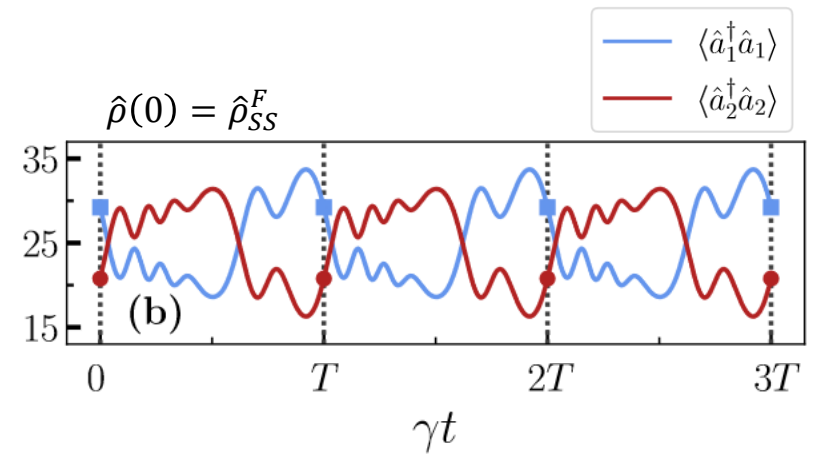
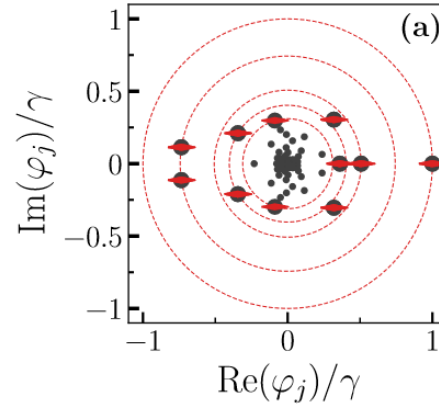
$$\Updownarrow$$

$$\mathcal{L}^F \hat{\rho}_j^F = \lambda_j^F \hat{\rho}_j^F$$

**Example model:** M. Hartmann et al, New J. Phys. (2017)

$$\hat{H}(t) = \frac{U}{2} \sum_{j=1}^2 \hat{a}_j^\dagger \hat{a}_j^\dagger \hat{a}_j \hat{a}_j - J(\hat{a}_1^\dagger \hat{a}_2 + \hat{a}_2^\dagger \hat{a}_1) + f(t)(\hat{a}_2^\dagger \hat{a}_2 - \hat{a}_1^\dagger \hat{a}_1)$$

$$\hat{L} = (\hat{a}_1^\dagger + \hat{a}_2^\dagger)(\hat{a}_1 - \hat{a}_2) = f_0 + f_1 \cos \omega t$$



## Conclusions

- Method to efficiently determine the low-lying eigenvalues and eigenmatrices of the evolution operator/Liouvillian;
- Retrieve these features through a (shorter) time evolution of the open system;
- Grants access to eigendecomposition of bigger system sizes.